

MODELING EXCESS HEAT IN THE FLEISCHMANN-PONS EXPERIMENT

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Over the past several years we have described models based on the linear coupling between two-level systems and an oscillator which are able to convert a large energy quantum into a large number of small energy quantum. At ICCF14 we presented a formula that can be used to calculate the phonon exchange matrix element for a nuclear reaction occurring in a lattice; in this presentation we report on our progress in the brute force computation of the matrix element based on realistic nuclear models. We consider also the development of semiclassical models which can approximate the dynamics of the quantum model. Finally, we describe the use of these models in a numerical simulation of the Fleischmann-Pons experiment. In these models a bottleneck occurs associated with helium diffusion out of the active region. We find that if the active region is thin (100 nm or less) that helium diffusion is fast, and excess heat production can occur in a steady state mode. If the active region is thick (500 nm or greater), then helium diffusion is slow and excess heat in these models occurs in bursts.



Modeling excess heat in the Fleischmann-Pons experiment

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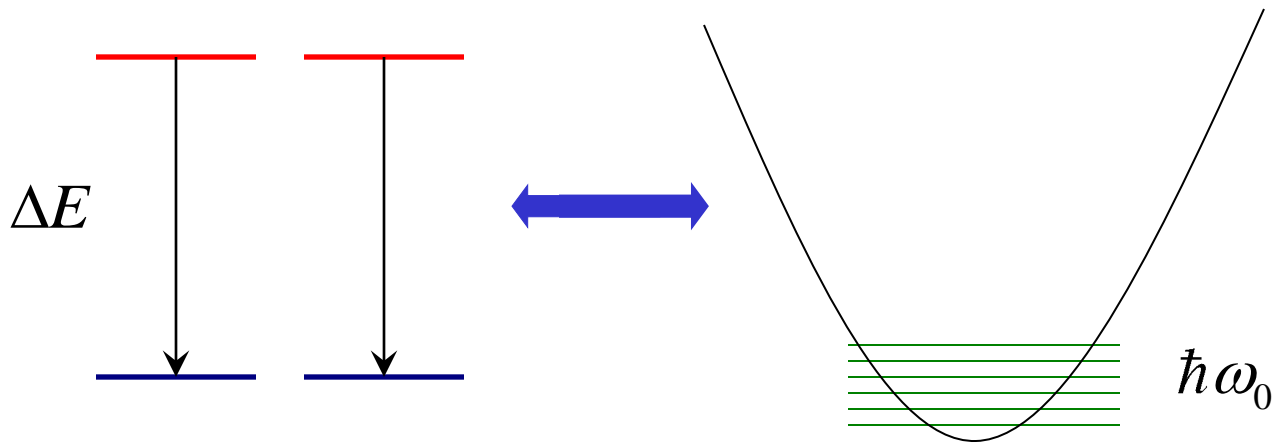
Theoretical problem

Although many more results available from experiment, we have enough so far to pose the key theory problem:

How to split up a large ΔE quantum into lots of small quanta?

The major implication of the Fleischmann-Pons experiment is that this is possible and occurs in energy production

Basic toy model



Two-level systems

Macroscopic
excited mode

$$\Delta E \gg \hbar\omega_0$$

Many-spin spin-boson model



C. Cohen-Tannoudji

$$\hat{H} = \Delta E \frac{\hat{S}_z}{\hbar} + \hbar \omega_0 \hat{a} \hat{a}^\dagger + V \frac{2S_x}{\hbar} (\hat{a} + \hat{a}^\dagger)$$

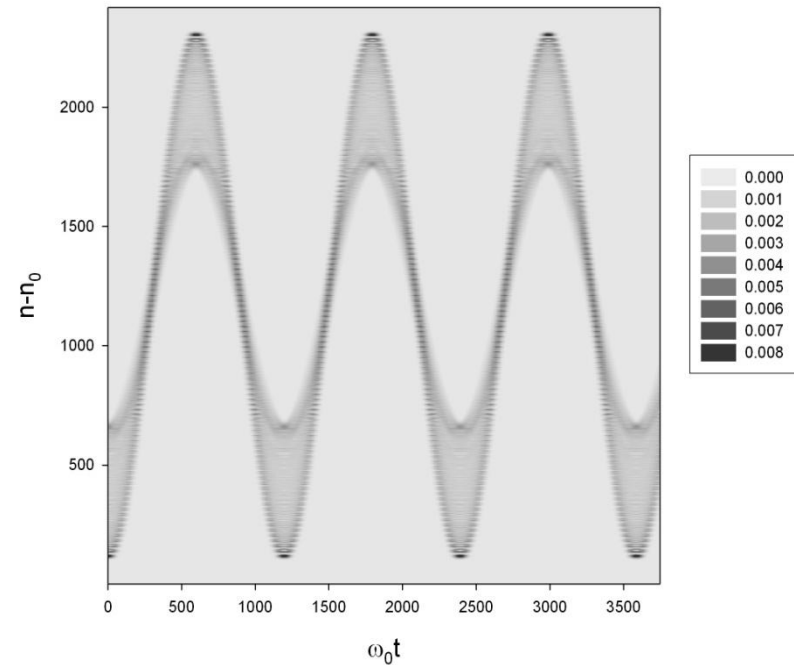
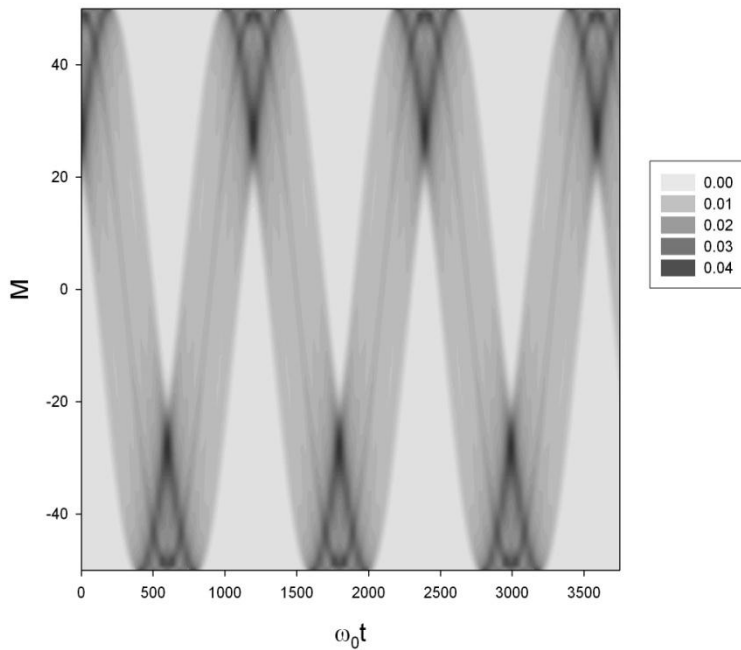
Two-level systems
energy

Harmonic oscillator
energy

Linear coupling
between two-level
systems and oscillator

Earlier versions of the model due to Bloch and Siegert (1940)

Coherent energy exchange



Numerical results for exchanging energy between 1700 oscillator quanta and 100 two-level systems

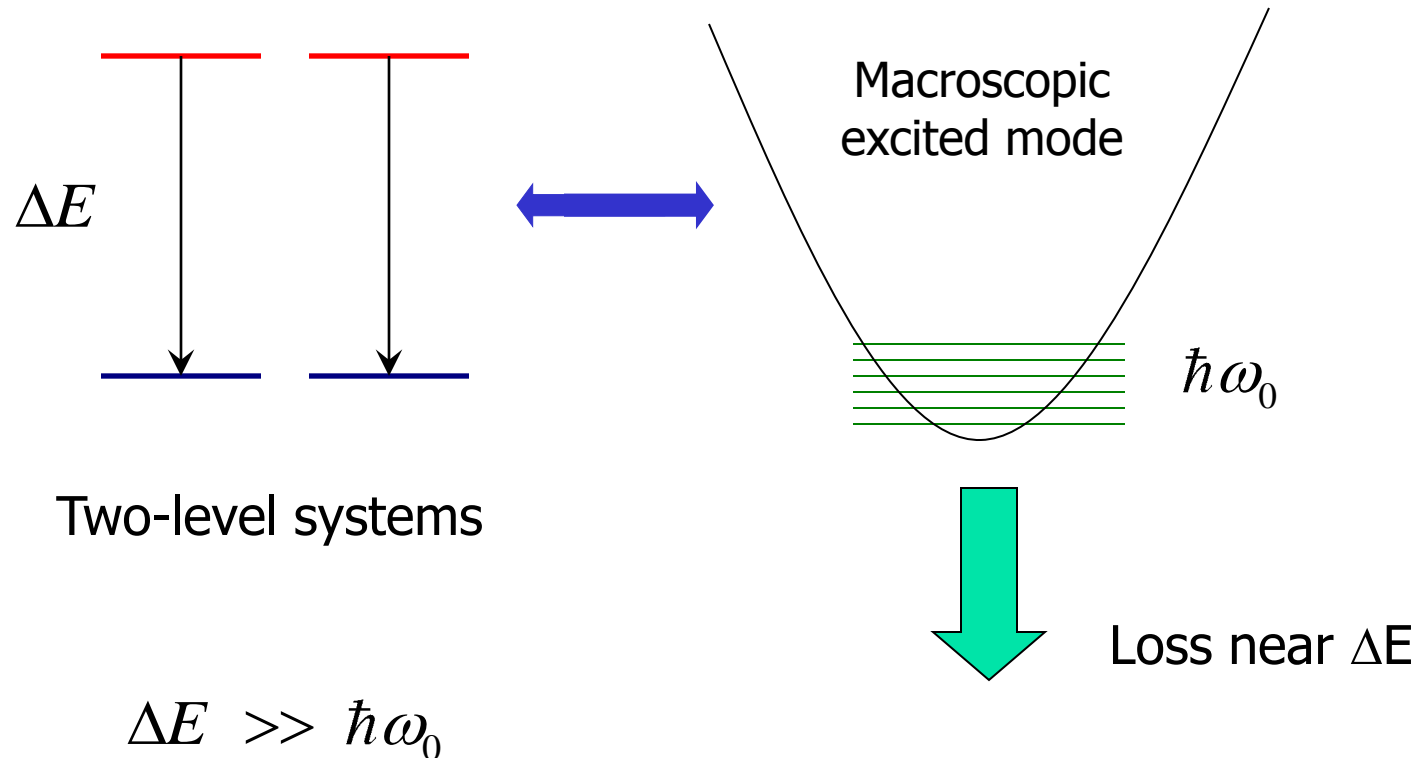


Thinking about toy model

Coherent multi-quantum energy exchange predicted by toy model

- Effect is weak
- Stringent resonance requirements
- Can exchange up to about 100 quanta coherently
- Exactly kind of model needed, except energy exchange effect is too weak


Improved toy model





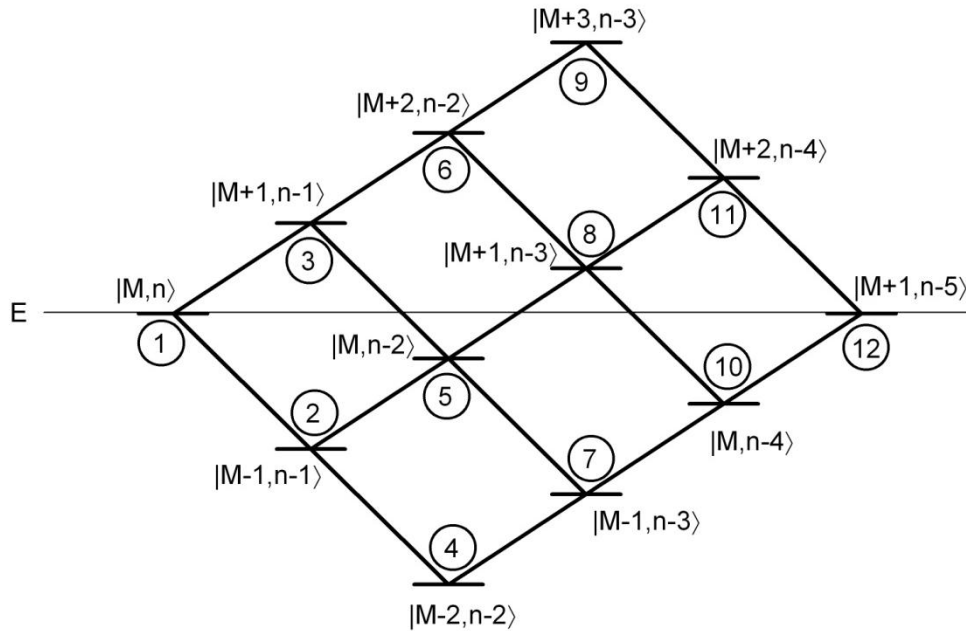
Lossy version of model

$$\hat{H} = \Delta E \frac{\hat{S}_z}{\hbar} + \hbar \omega_0 \hat{a} \hat{a}^\dagger + V \frac{2S_x}{\hbar} (\hat{a} + \hat{a}^\dagger) - i \frac{\hbar}{2} \Gamma(E)$$



Loss term, which allows the system to decay when a large energy quantum is available

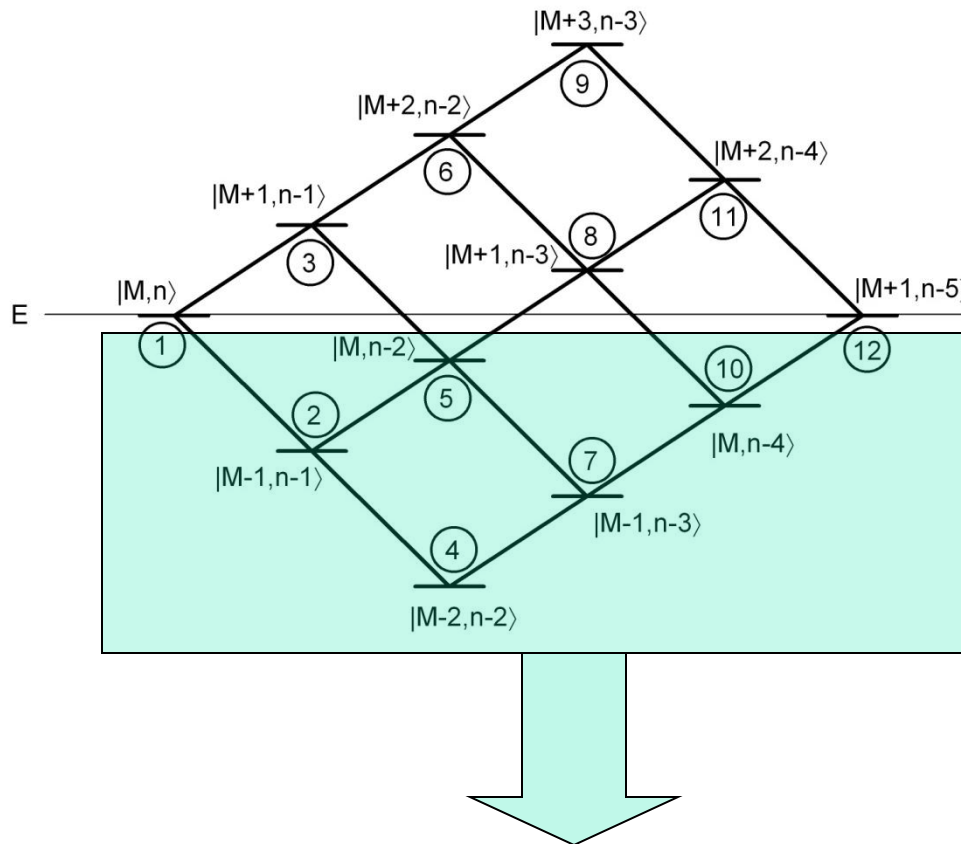
Perturbation theory



Many paths from initial to final state, with interference between upper and lower paths

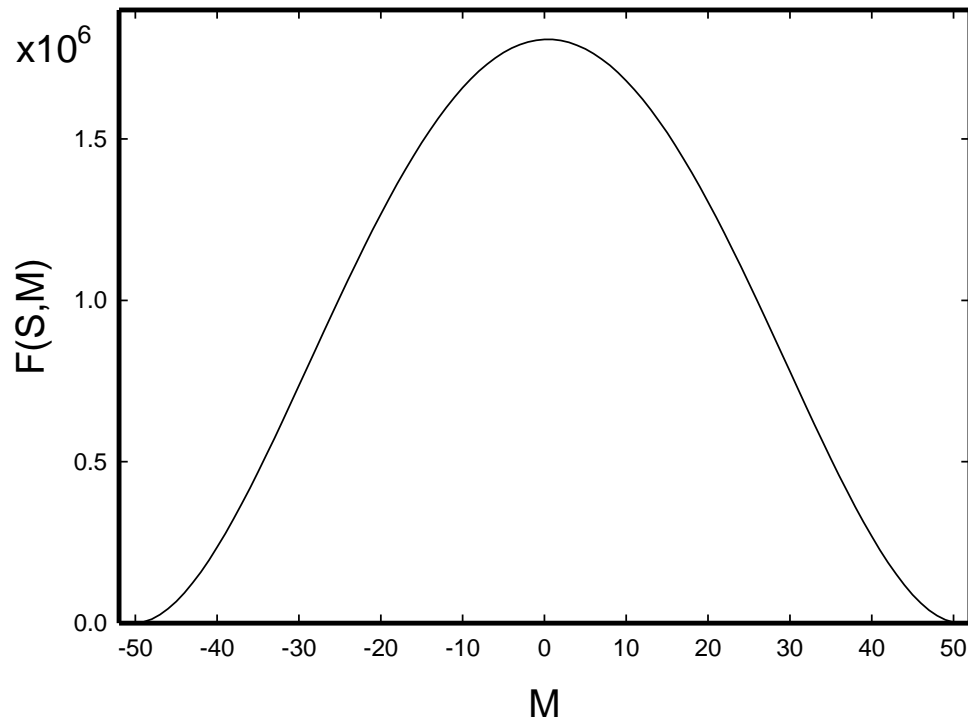
Finite basis approximation for $|n\rangle \otimes |M\rangle \rightarrow |n-5\rangle \otimes |M+1\rangle$

Perturbation theory



Loss channels available for off-resonant states with energy excess, which spoils the destructive interference

Enhancement due to loss



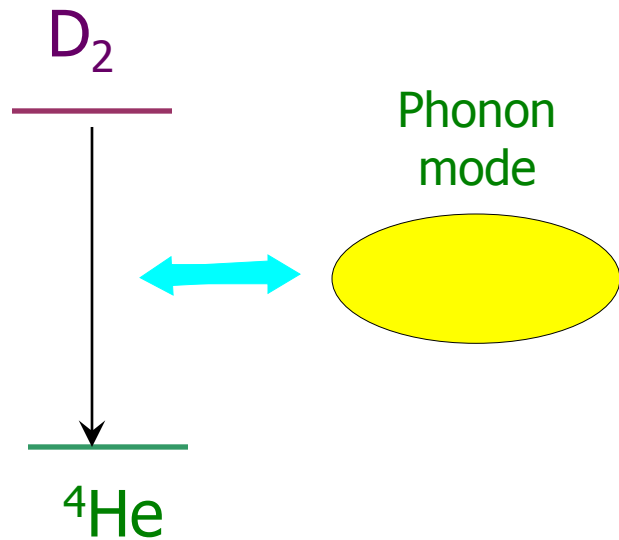
$$\left[V_{1,12}(E) \right]_{\Gamma=\infty} = \left[V_{1,12}(E) \right]_{\Gamma=0} F(S, M)$$



Lossy version of model

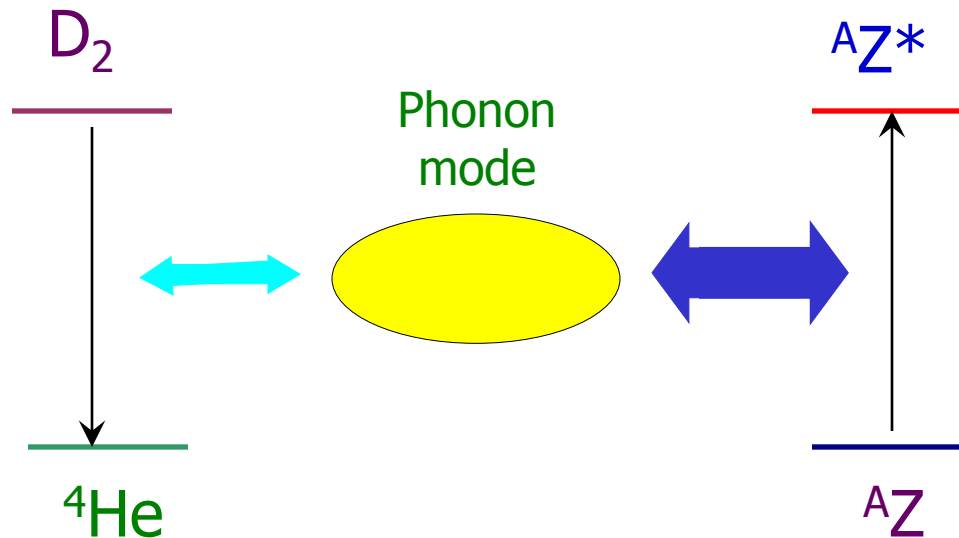
- Loss spoils the destructive interference
- Coherent energy exchange rates increased by orders of magnitude
- Much stronger effect
- Model capable of converting 24 MeV to atomic scale quanta

Thinking about PdD



Unfortunately, coupling is too weak because of Coulomb repulsion

Excitation transfer



Indirect evidence from experiment implicates $\text{AZ} = {}^4\text{He}$, and theory and experiment suggest that AZ^* is a localized two-deuteron state

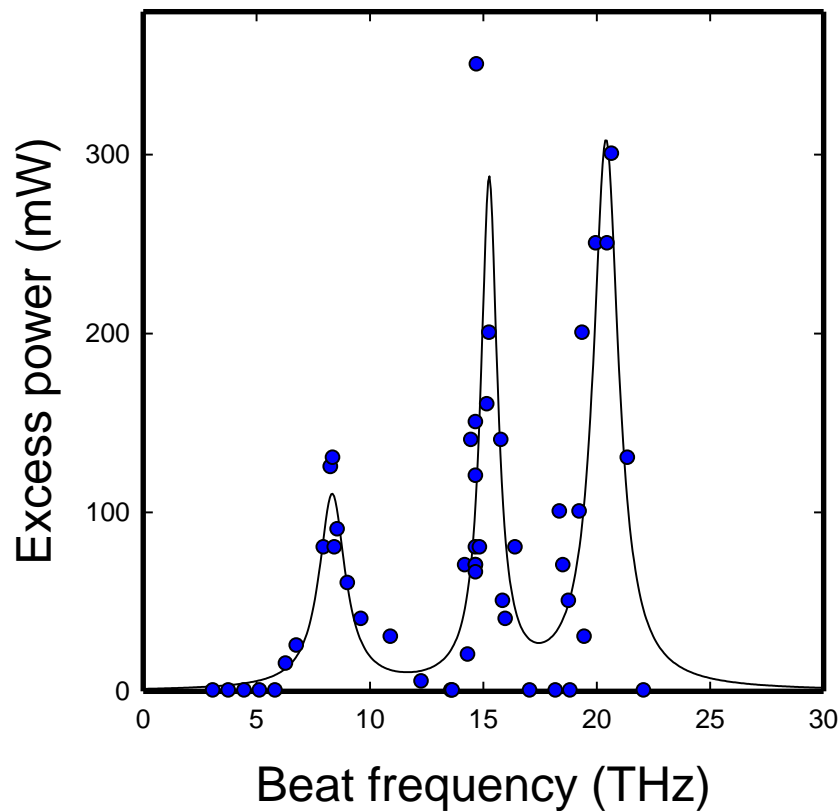


Basic model

$$\begin{aligned}\hat{H} = & \Delta E_1 \frac{\hat{S}_z^{(1)}}{\hbar} + \Delta E_2 \frac{\hat{S}_z^{(2)}}{\hbar} + \hbar \omega_0 \hat{a} \hat{a}^\dagger - i \frac{\hbar}{2} \Gamma(E) \\ & + V_1 e^{-G} \frac{2S_x^{(1)}}{\hbar} (\hat{a} + \hat{a}^\dagger) + V_2 \frac{2S_x^{(2)}}{\hbar} (\hat{a} + \hat{a}^\dagger)\end{aligned}$$

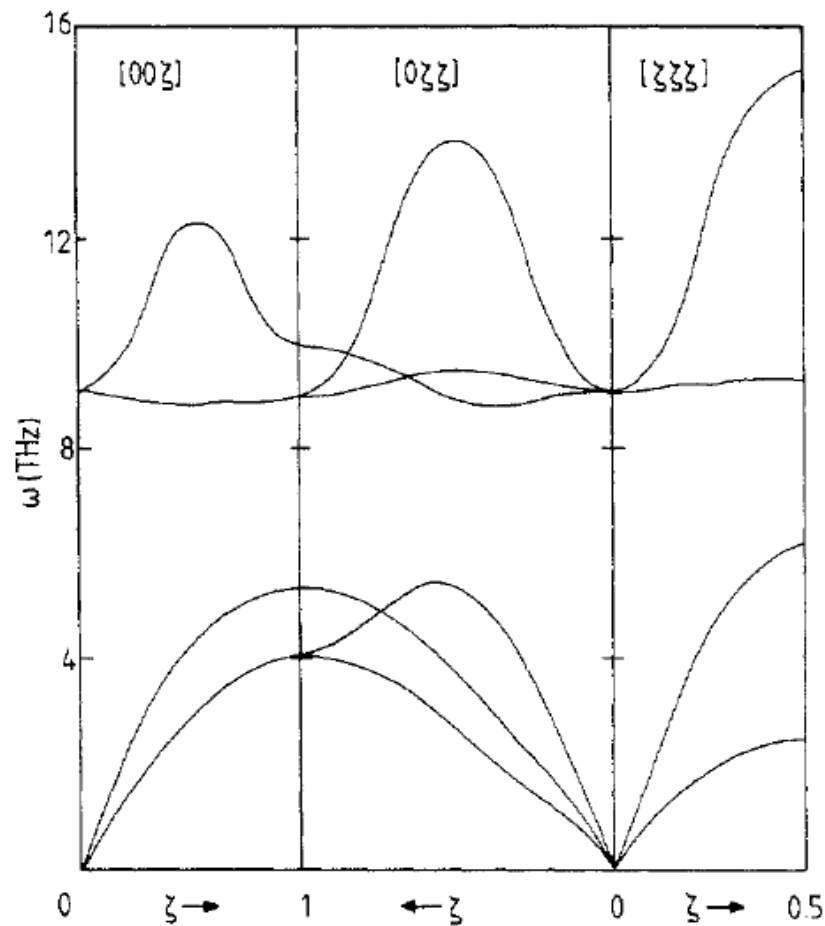
This kind of model is first one relevant to experiment

What oscillator modes?



Frequency (THz)	width (THz)
8.3	0.70
15.3	0.44
20.4	0.68

Dispersion curve for PdD



L E Sansores et al
J Phys C **15** 6907 (1982)



Strong-coupling limit

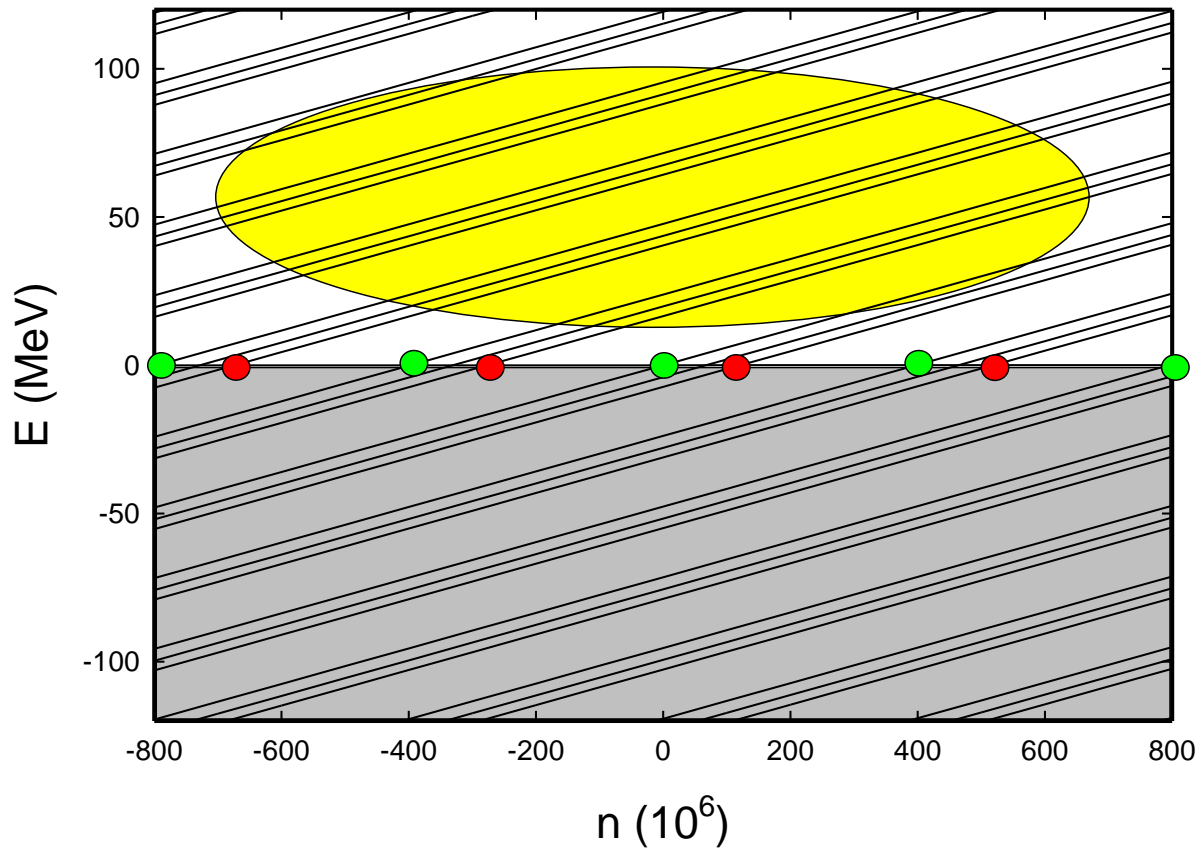
When the coupling between the receiver-side two-level systems and oscillator is strong, then the problem simplifies

$$\Gamma \rightarrow \frac{\hbar\omega_0}{\Delta E(g)} \left| \frac{\langle S, M, n + \Delta n | \hat{H} | S, M + 1, n \rangle}{\hbar} \right|$$

When the excitation transfer step is the bottleneck, then

$$\Gamma = \frac{V_1\sqrt{n}}{\hbar} \left(\frac{\hbar\omega_0}{\Delta E} \right) e^{-G} \sqrt{(S+M)(S-M)}$$

Occupation of virtual levels





Conclusions so far

- Can model the effect
- Can see the energy exchange with the lattice
- Can see excitation transfer
- Can get rates for both
- Agreement with experiment if screening energy $U_e = 150$ eV



Trying out simplified version

$$\frac{d}{dt} N_{D2} + \frac{N_{D2} - N_{D2}^0}{\tau_{D2}} = -\Gamma_0 \sqrt{N_{D2} N_{He}} \Theta(n - n_{thresh})$$

$$\frac{d}{dt} N_{He} + \frac{N_{He} - N_{He}^0}{\tau_{He}} = \Gamma_0 \sqrt{N_{D2} N_{He}} \Theta(n - n_{thresh})$$

$$\frac{d}{dt} n + \frac{n - n_0}{\tau_p} = \gamma_J + \frac{\Delta E}{\hbar \omega_0} \Gamma_0 \sqrt{N_{D2} N_{He}} \Theta(n - n_{thresh})$$



Example: fast He diffusion

Active region:

$$A = 1 \text{ cm}^2$$

$$\Delta r = 100 \text{ nm}$$

D₂ parameters:

$$f[\text{vacancy}] = 0.25$$

$$f[\text{D}_2] = 0.005$$

$$N[\text{D}_2] = 1.8 \times 10^{15}$$

$$\tau_{\text{D}_2} = 2 \times 10^{-8} \text{ sec}$$

⁴He parameters:

$$D_{\text{He}} = 1.3 \times 10^{-14} \text{ cm}^2/\text{sec}$$

$$\tau_{\text{He}} = \Delta r^2 / D_{\text{He}} = 2.1 \text{ hr}$$

Phonon mode:

$$f_0 = 8.3 \text{ THz}$$

$$Q = 20$$

Deuterium flux:

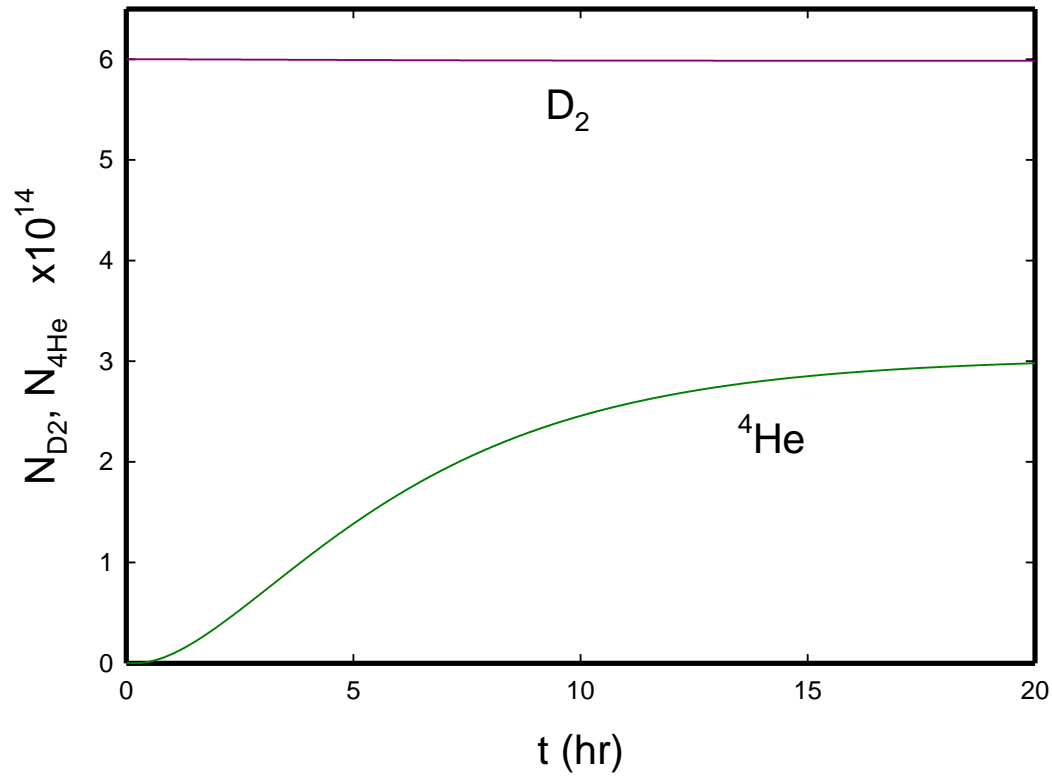
$$P_{\text{flux}} = 1 \text{ Watt/cm}^3$$

$$n_{\text{thresh}} = 100$$

Basic reaction rate:

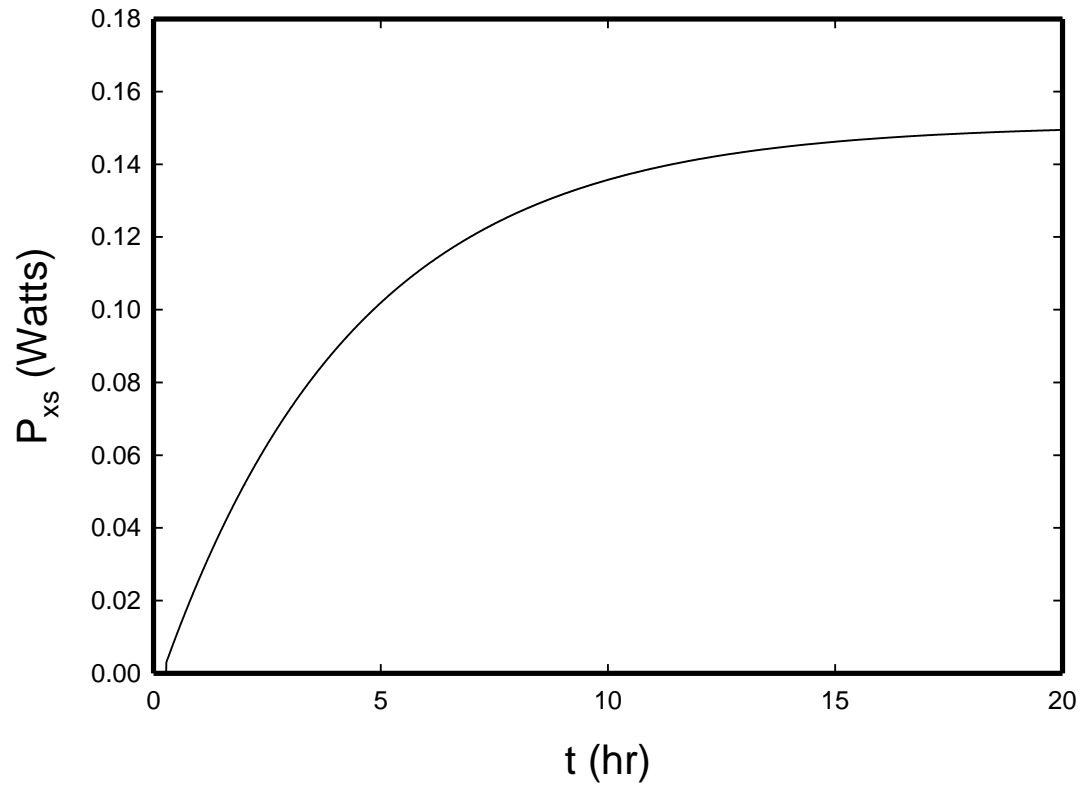
$$\Gamma_0 = 1/(3 \text{ hr})$$

Evolution of dideuterium, ^4He

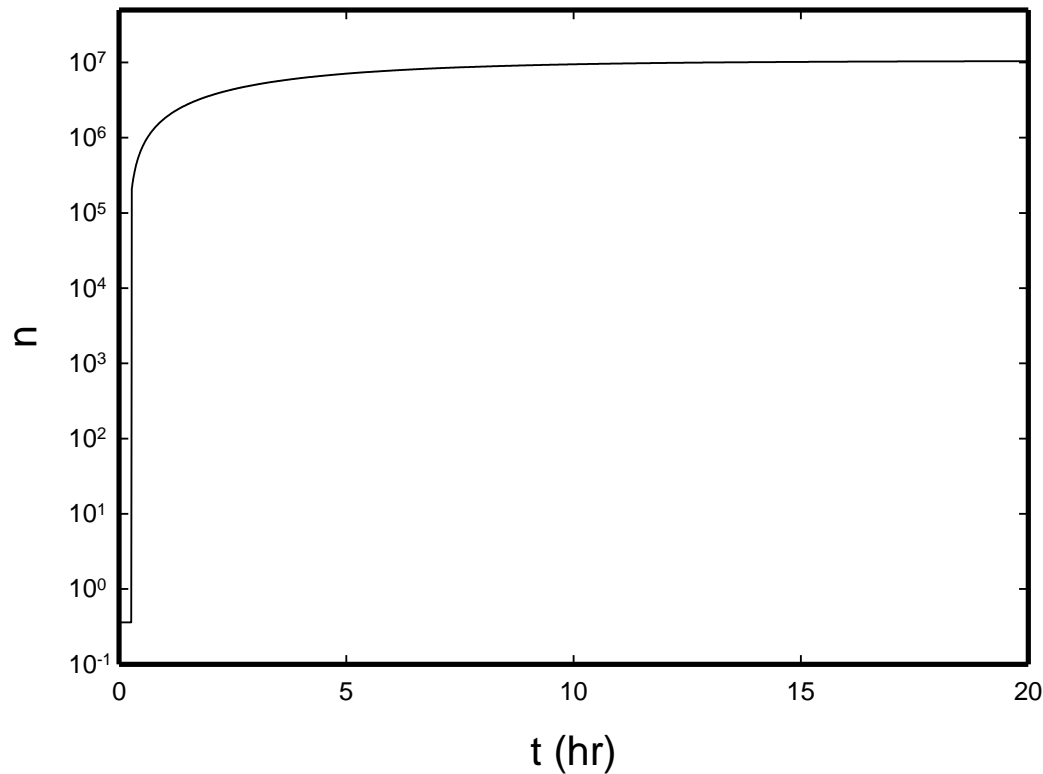




Excess power



Number of phonons



Thermal: 0.36 Flux generated (1 W/cm^3): 700 P_{XS} generated: 10^7



Addressing the full problem

Start out with full problem

$$E\Psi = H\Psi$$

Then implement picture and approximation through construction of channels

$$\Psi = \sum_j \psi_j \Phi_j$$

Get coupled-channel equations

$$\langle \Phi_i | E\Psi = \langle \Phi_i | H\Psi$$



Coupled-channel equations

Coupled-channel equations that result

$$E\psi_i = H_{ii}\psi_i + \sum_{j \neq i} H_{ij}\psi_j$$

Can put whatever physics that one likes into the channels.
Best place to start is with

$$H = H[\text{nucleons}] + H[\text{electrons}] + V[\text{strong force}] + V[\text{Coulomb}]$$



Coupling

Terms that couple from one channel to another:

$$H_{ij} = \langle \Phi_i | H | \Phi_j \rangle$$

We focus on strong force terms, although others present

$$H_{ij} = V_n e^{-iS_D}$$

Strong force Lattice change

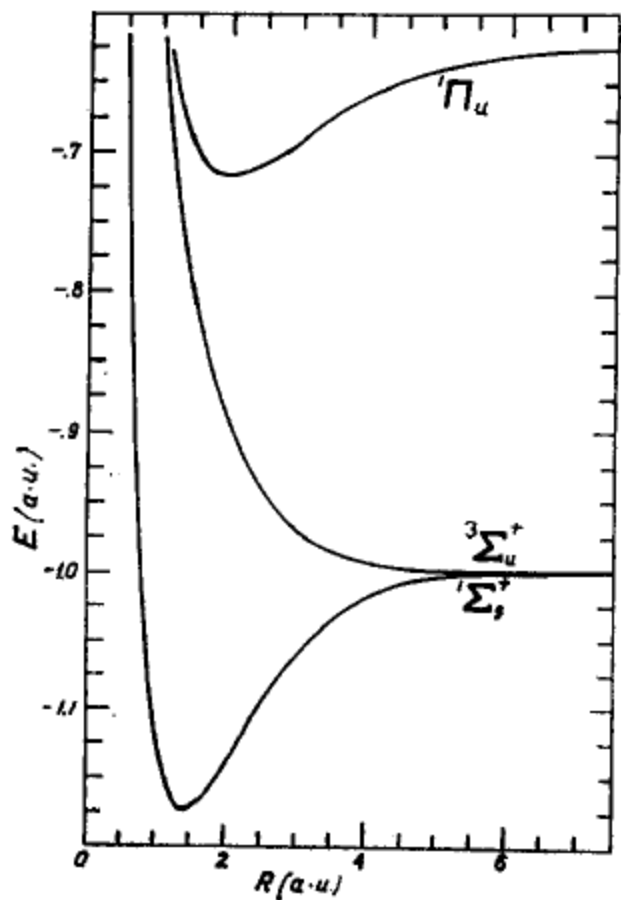


Where is the D_2 ?

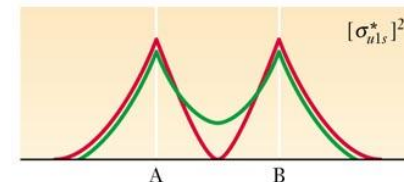
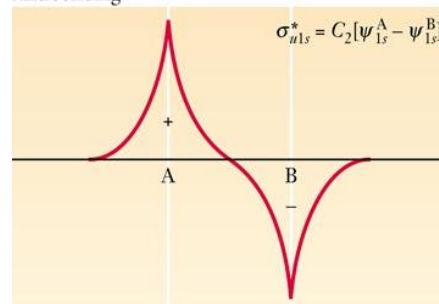
Molecular D_2 does not form in bulk PdD

- Issue is electron density
- Computation of D_2 in electron gas leads to occupation of antibonding states
- The electron density in PdD is too high
- If you want D_2 , you have to lower the electron density

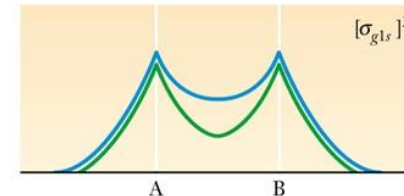
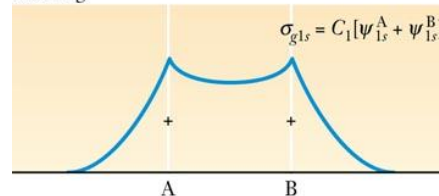
Bonding and anti-bonding in H₂



Antibonding

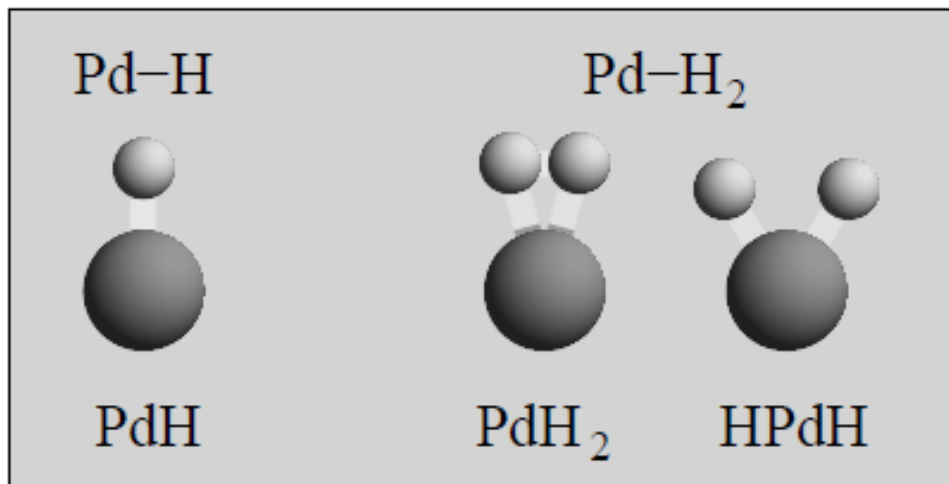


Bonding



W. Kolos and L. Wolniewicz,
J Chem Phys **43** 2429 (1965)

Pd-H₂ with σ -bonding



In Pd-H₂

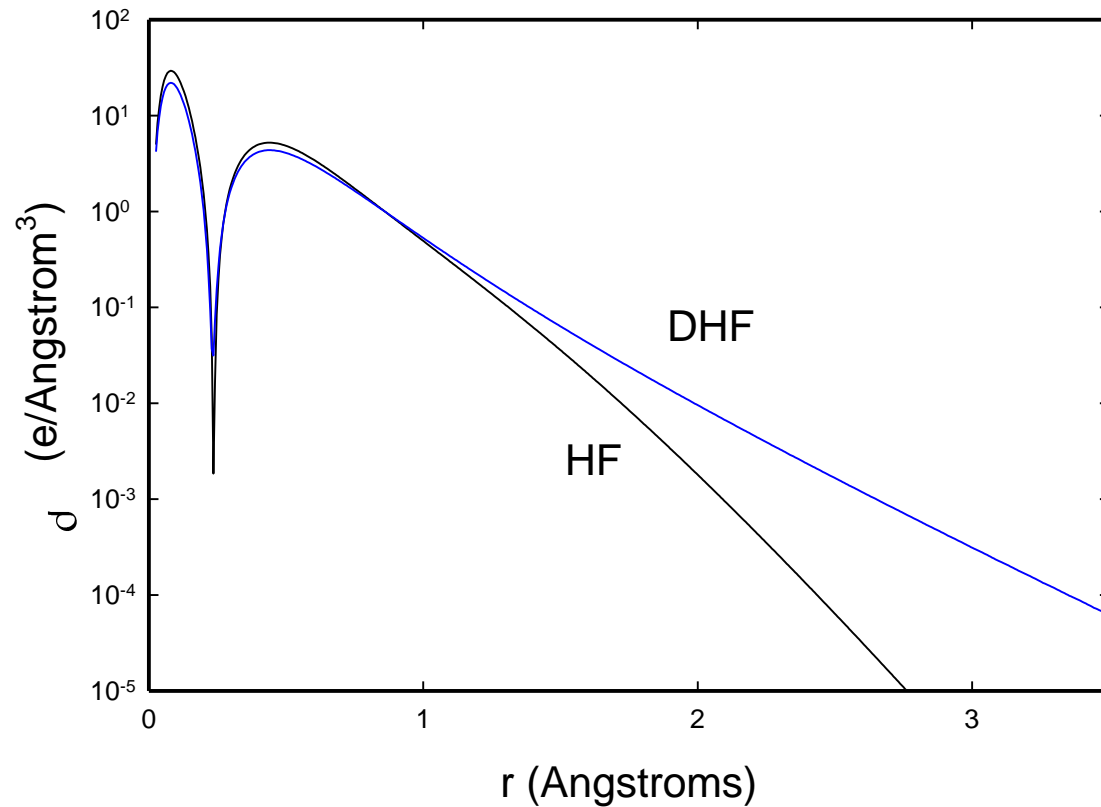
$d[\text{Pd-H}] = 1.67 \text{ \AA}$

$d[\text{H-H}] = 0.81 \text{ \AA}$

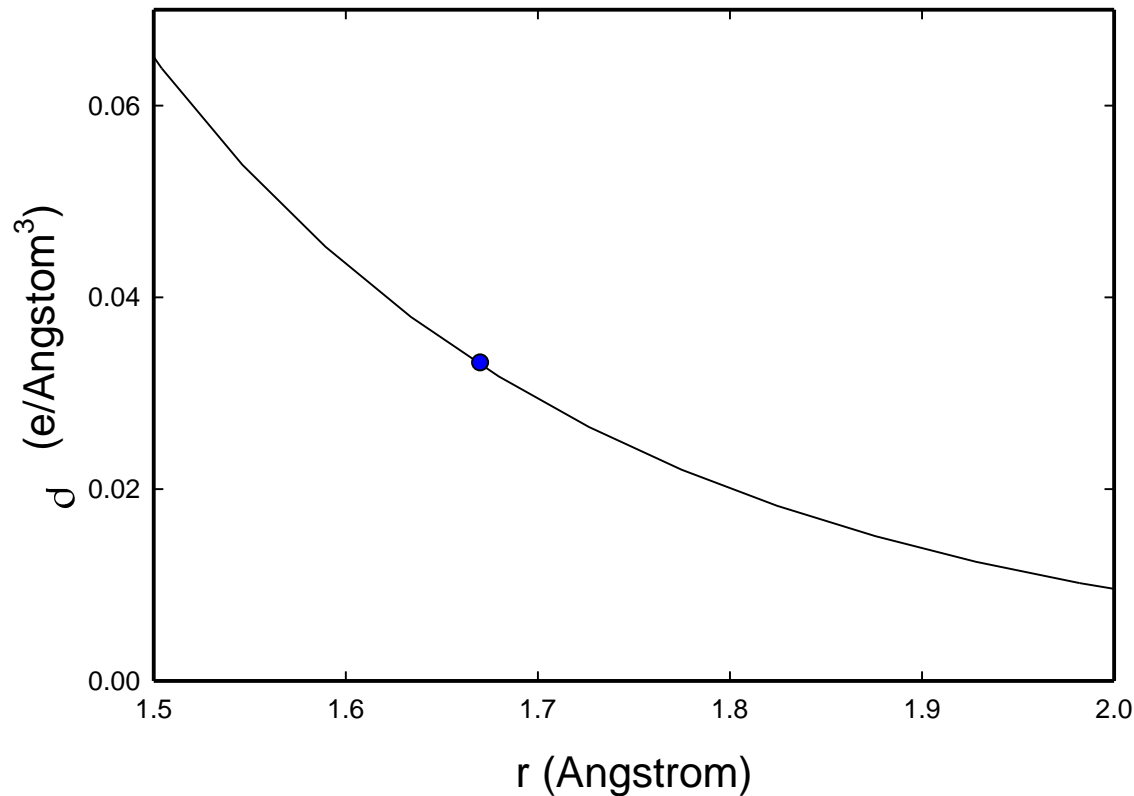
σ -bonded Pd-H₂ is the ground state of the three-atom system.

It is a combination of $(4d)^{10} \text{ } ^1S_0$ Pd and $(1\sigma)^2 \text{ } ^1S_0$ H₂

Electron density of Pd (4d)¹⁰

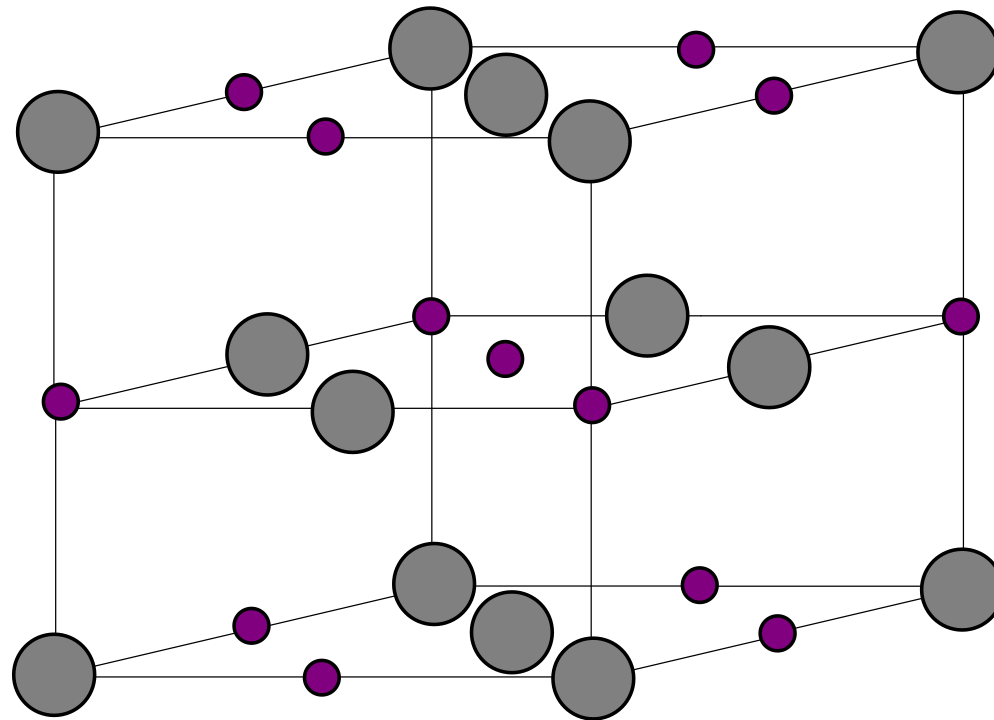


Electron density at Pd-H distance

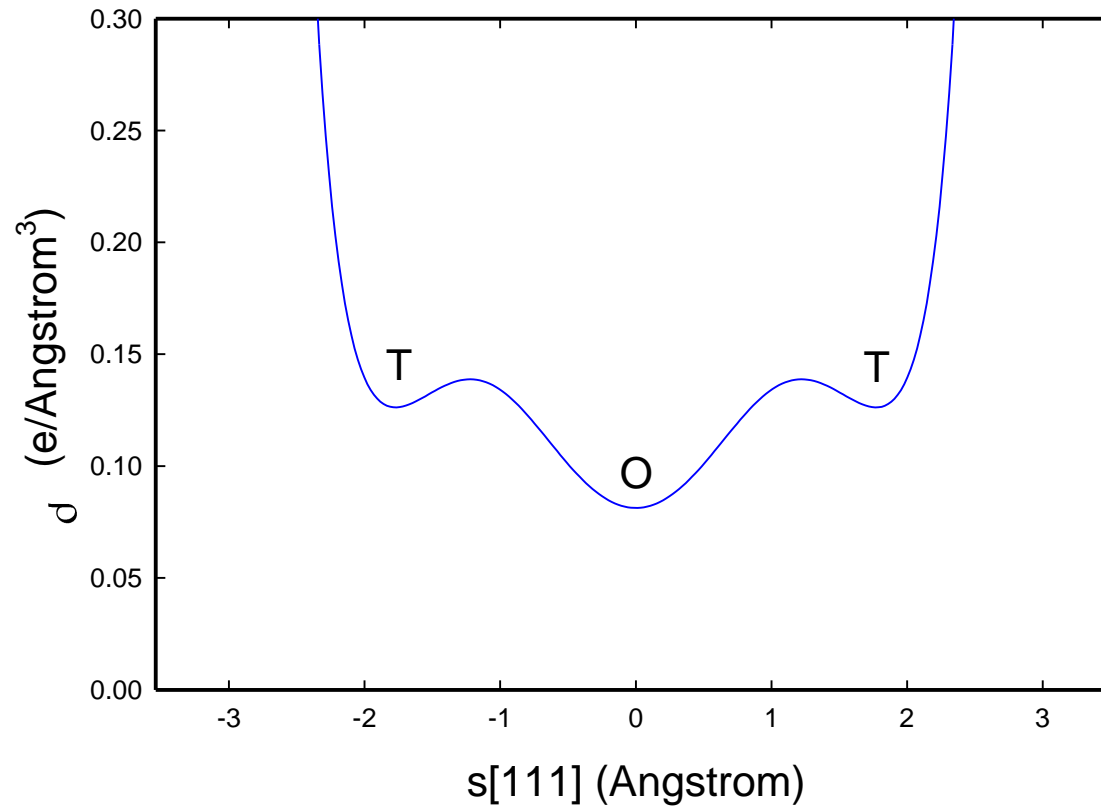


Pd-H distance in Pd-H₂ is 1.67 Angstroms, and electron density is 0.033

PdD lattice structure (fcc)



Electron density due to Pd around octahedral site





Cannot form D_2 at O-site

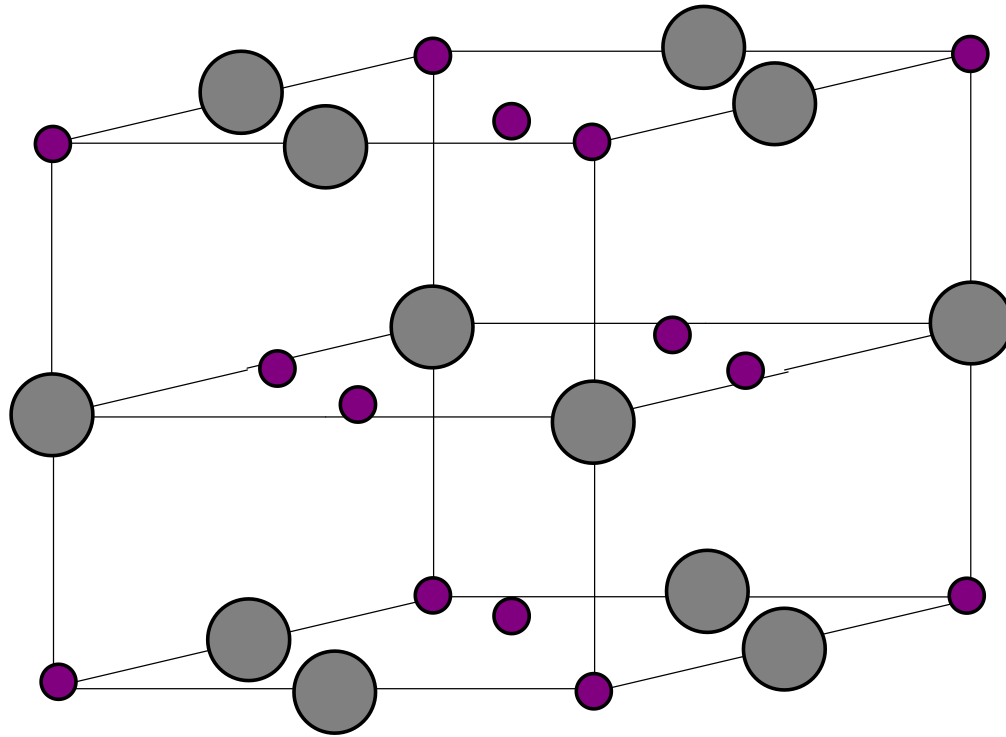
H_2 binds with Pd at 1.67 Angstroms Pd-H separation

Electron density at 1.67 Angstroms is $0.33 \text{ e}/\text{Angstrom}^3$

Electron density at O-site in PdD is $0.081 \text{ e}/\text{Angstrom}^3$

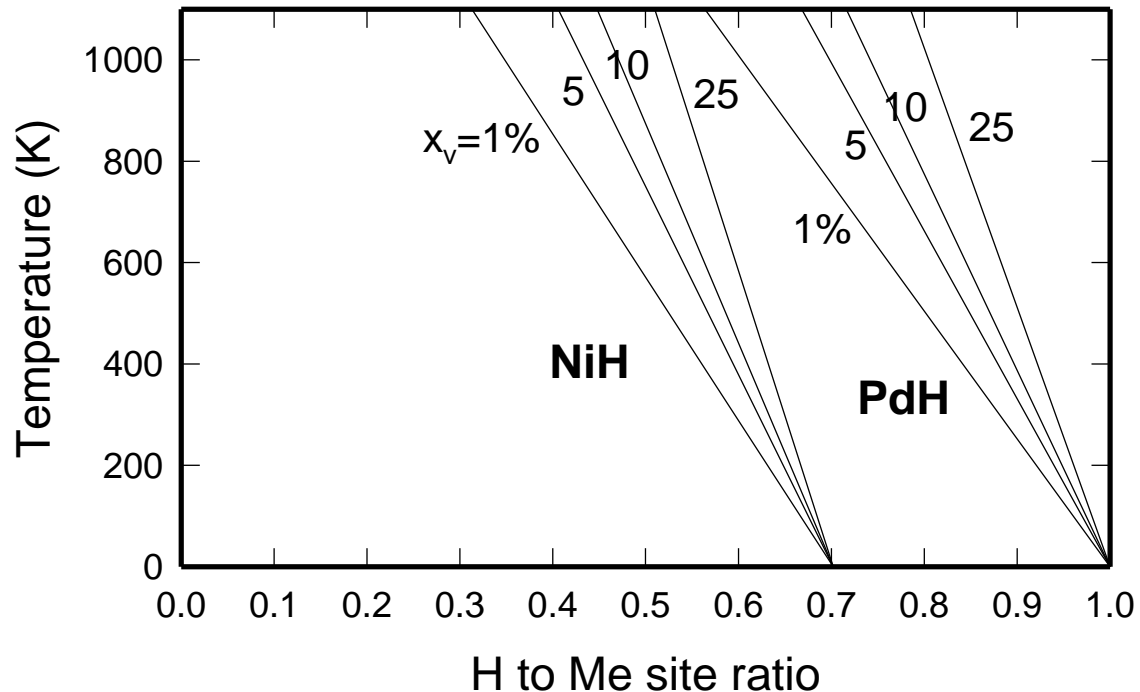
Anti-bonding orbitals occupied

PdD Host lattice vacancy



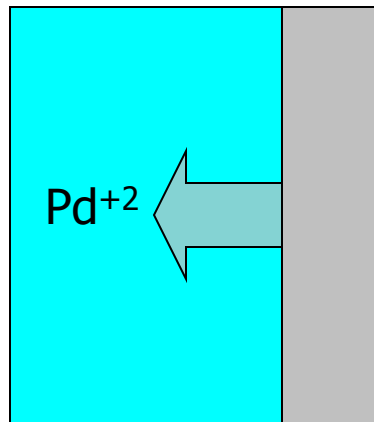
Deuterium atoms relax toward host vacancy

Vacancies in host lattice

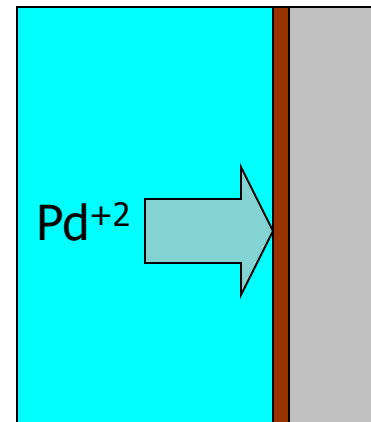


Vacancies in host metal lattice are thermodynamically favored at high loading

Codeposition



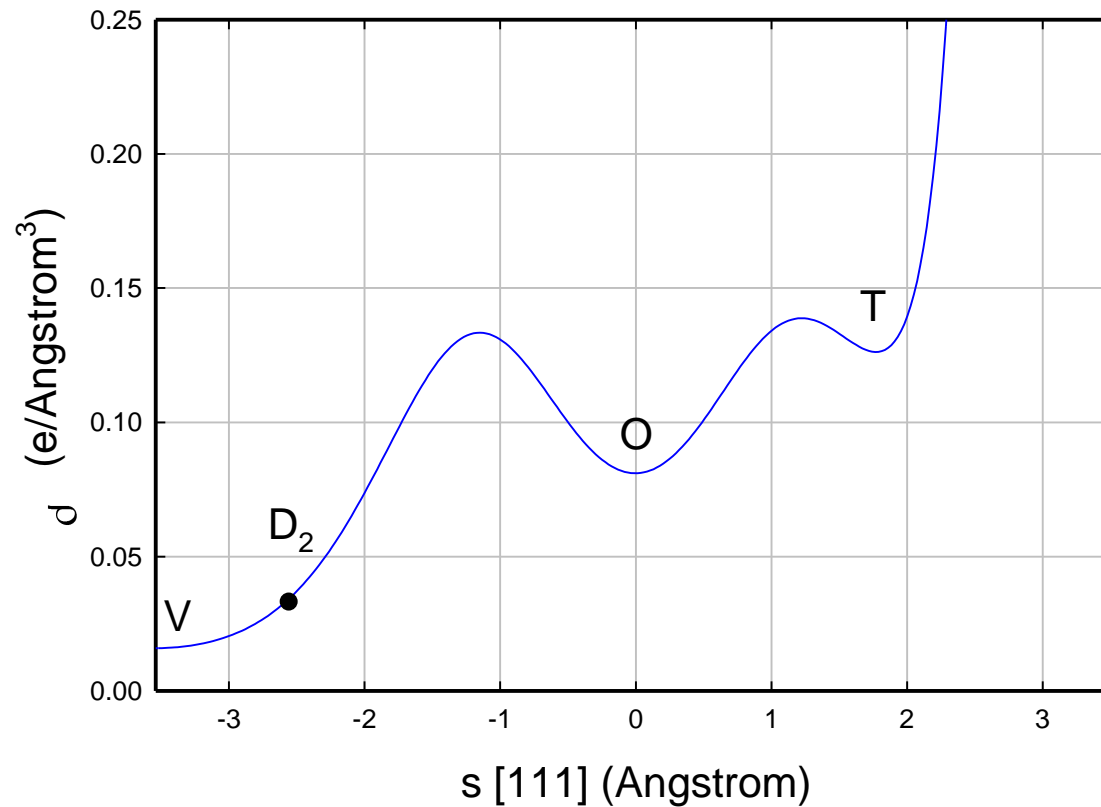
Anodic current



Cathodic current

Conjecture that a small amount of Pd is stripped off during anodic current cycles, and then codeposited during subsequent cathodic loading [most of the Pd in solution is $\text{Pd}(\text{OH})_4^{-2}$, Mountain and Wood (1988)]

Electron density with vacancy



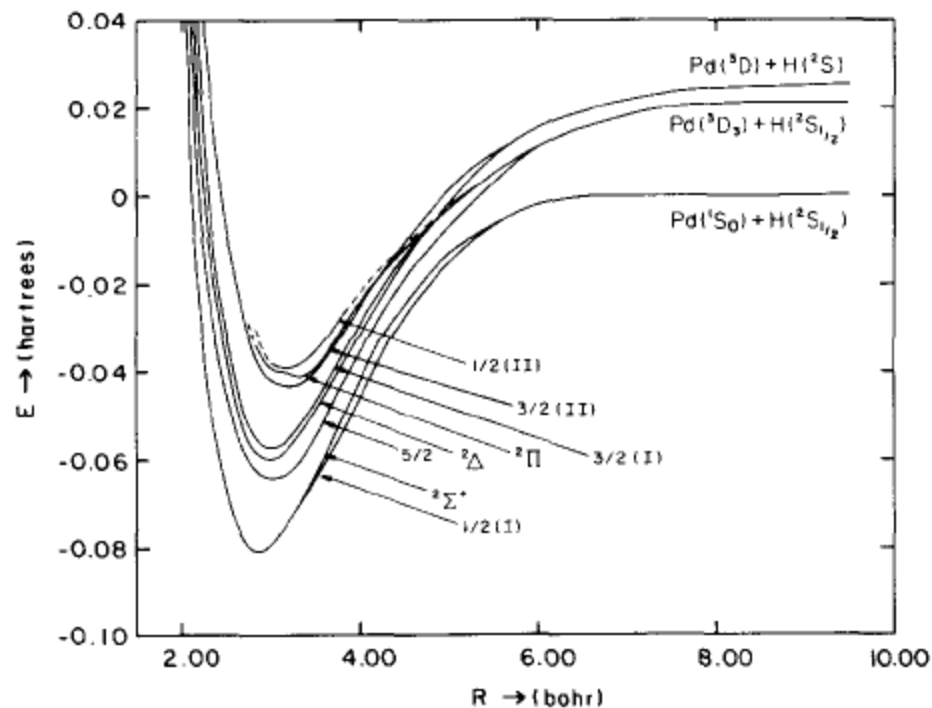


Electron density seems low enough

Superposition of atomic electron densities leads to a model electron density of $0.016 \text{ e}/\text{Angstrom}^3$

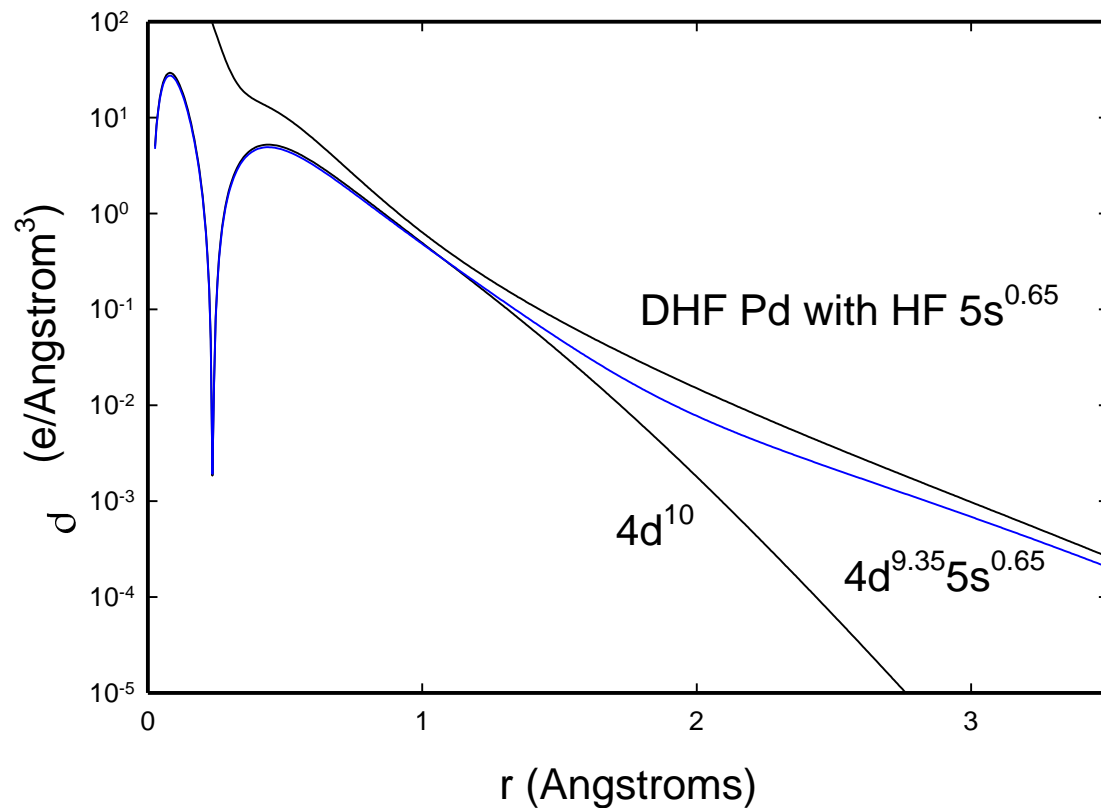
- Model electron density is 2x lower than for Pd-H₂
- Would expect D₂ formation near vacancy
- Would expect relevant literature

Check with PdH



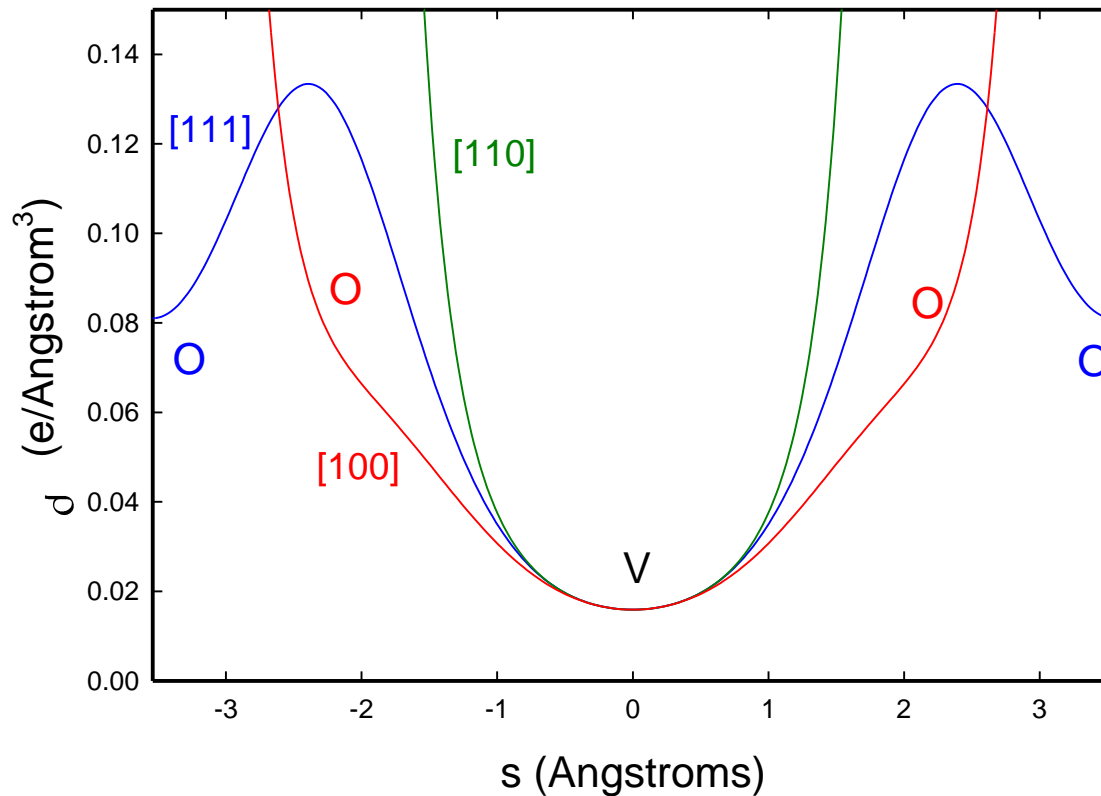
For ground state
 $d[\text{Pd-H}] = 1.53 \text{ \AA}$

Look at Pd density at 1.53 A



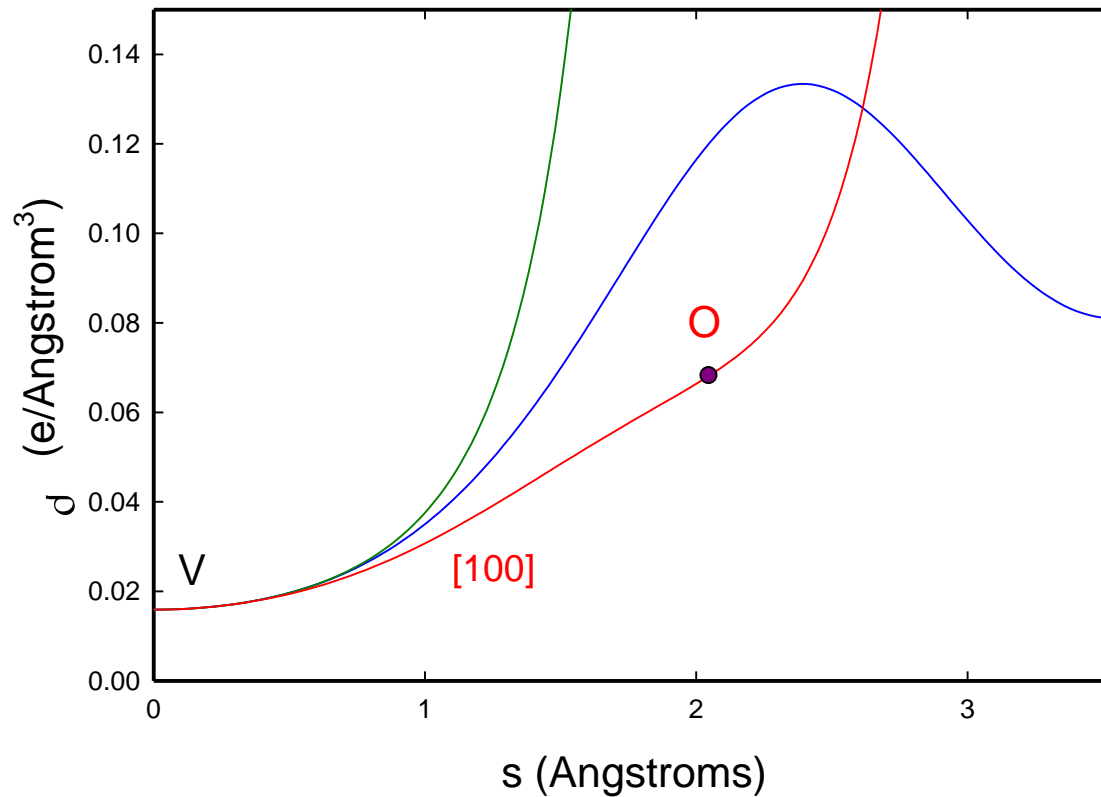
Model electron density at 1.53 A is $0.0686 \text{ e}/\text{Angstrom}^3$

Electron density around vacancy

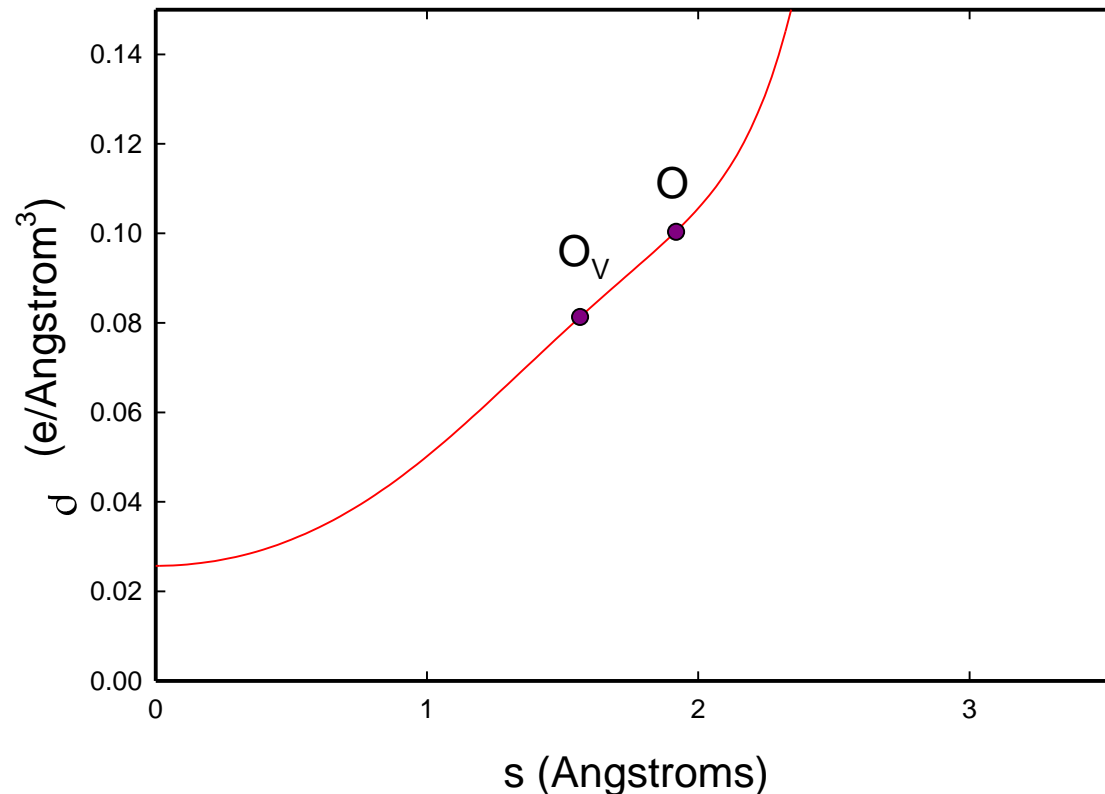


Nearest O-point no longer a minimum in electron density

Model electron density just right now for H



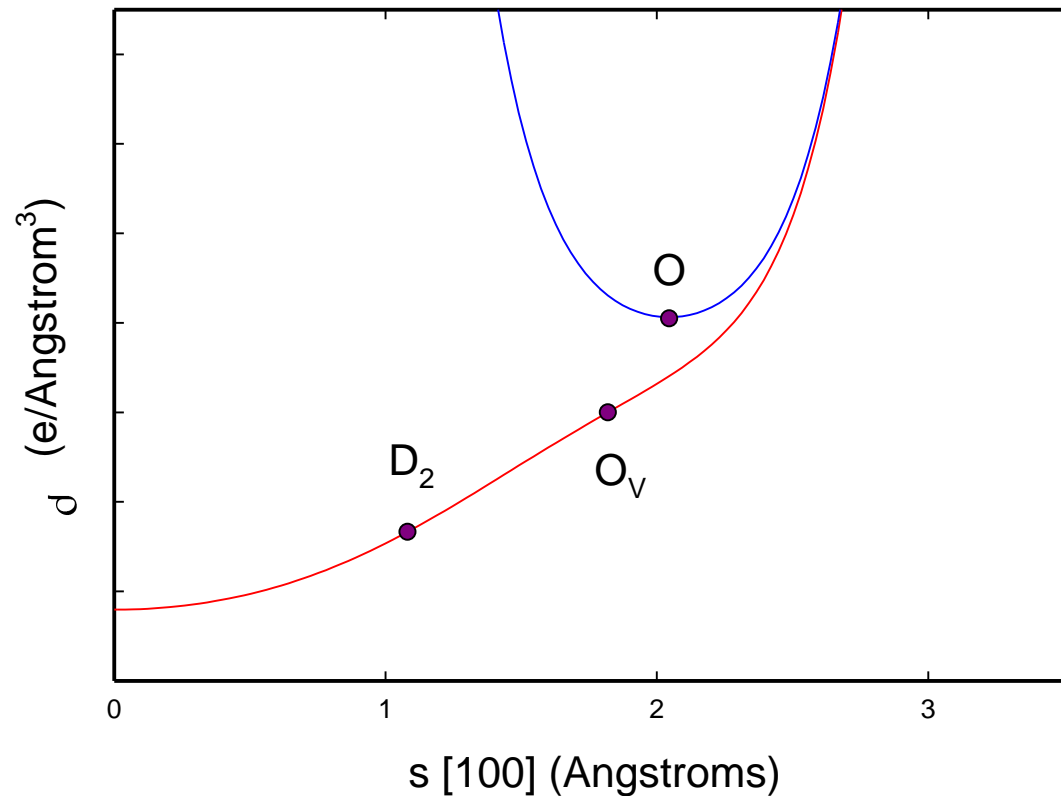
Compare with Velikova et al (2009) DFT for Pd



Velikova et al,
Phys Rev B **80**
024101 (2009)

Velikova shift corresponds to 0.081 e/Angstrom³ , close to PdH 0.069 e/Angstrom³

Summary (need VASP calculation!)



Expect about 0.4 Å shift of D, and about 1 Å shift for D₂ location



Summary and conclusions

- Fleischmann-Pons experiment points to new kind of physical process where nuclear energy generated with no energetic reaction products
- Spin-boson model provides analog which can convert a big quantum to a large number of small quanta, but effect is weak
- Lossy spin-boson model can convert a large number a big quantum to a large number of small quanta, and effect is large



More conclusions

- We can construct coupled-channel equations systematically to implement lossy spin-boson type scheme in real physical system
- Detailed modeling now under way